

Highly Stable Perovskite Photodetector Based on Vapor-Processed Micrometer-Scale CsPbBr₃ Microplatelets

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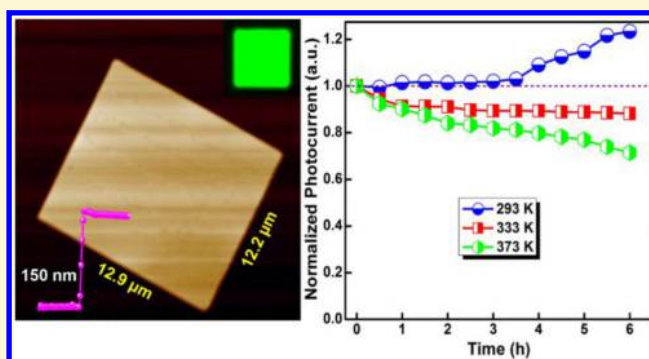
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Supporting Information

ABSTRACT: Recently, halide perovskites have attracted tremendous attention because of their great abilities in large-scale and cost-effective manufacturing of optoelectronic devices. Here, a novel photodetector configuration was proposed by employing vapor-processed micrometer-scale inorganic CsPbBr₃ microplatelets as the light absorber. Temperature-dependent steady-state and time-resolved photoluminescence spectra were first recorded to study the emission mechanisms and carrier recombination dynamics of the CsPbBr₃ microplatelets. Furthermore, a photoconductive detector was prepared, and the device exhibits good performances with a high on/off photocurrent ratio of 4.6×10^3 , a responsivity of ~ 1.33 A/W, and a specific detectivity of 0.86×10^{12} jones. Additionally, temperature-dependent current–voltage and current–time characteristics of the photodetector were studied to assess the thermal effects on its photodetection ability. In particular, the unencapsulated photodetector demonstrates a prominent stability over the long-term temperature endurance measure in ambient air. Even operated at 373 K, the photodetector can operate properly, showing a high temperature resistance. Moreover, the device performance can almost be retained even with a 7 month storage in air. The experimental results suggest that the CsPbBr₃ microplatelets can serve as a good candidate for the fabrication of high-performance photodetectors compatible with practical applications.



INTRODUCTION

Recently, metal halide perovskites have triggered an increasingly intensive interest owing to their remarkable characteristics, including direct bandgap, tunable emission wavelength, high light absorption coefficient, simple processing technique, and high and well-balanced carrier transport ability.^{1–10} These unique features render such perovskite materials promising candidates for optoelectronic device applications such as photovoltaic cells,^{11,12} luminescent devices,^{13–16} lasing,^{17–19} and photodetectors.^{4,20–24} However, the study of photodetectors is still short of important breakthroughs that can promote perovskite-based photodetectors into practical applications, which is mainly because of the difficulty in preparing high-quality perovskite light absorbers with a good compactness and a good crystallinity and also the undesired hypersensitivity of perovskites in air with moisture. It is well-known that when organic–inorganic hybrid perovskites are exposed to air ambient, they suffer rapid deterioration. However, inorganic perovskites CsPbX₃ show remarkable

stability except for all the superior characteristics inherited from the organic–inorganic perovskites, and many reports have argued that CsPbX₃ has a high thermal stability until its melting at ~ 500 °C.²⁵ Thus, as the alternatives, inorganic perovskites are beginning to draw much scientific attention in reliable photodetector fabrication.^{26–31} For instance, Chen et al. demonstrated a simple photodetector using horizontal single-crystal CsPbX₃ nanowires with controlled crystallographic orientations as the light absorber, showing an on/off photocurrent ratio of 10^3 and a response time of <0.1 s.²⁶ Lu et al. reported a CsPbBr₃ photodetector integrated with Si nanowires with a responsivity of 54 mA/W and a response time of 0.48/1.03 ms.²² Zeng and co-workers demonstrated a planar device configuration constructed with CsPbBr₃ nanocrystals with a preferred orientation, and the photoresponsivity

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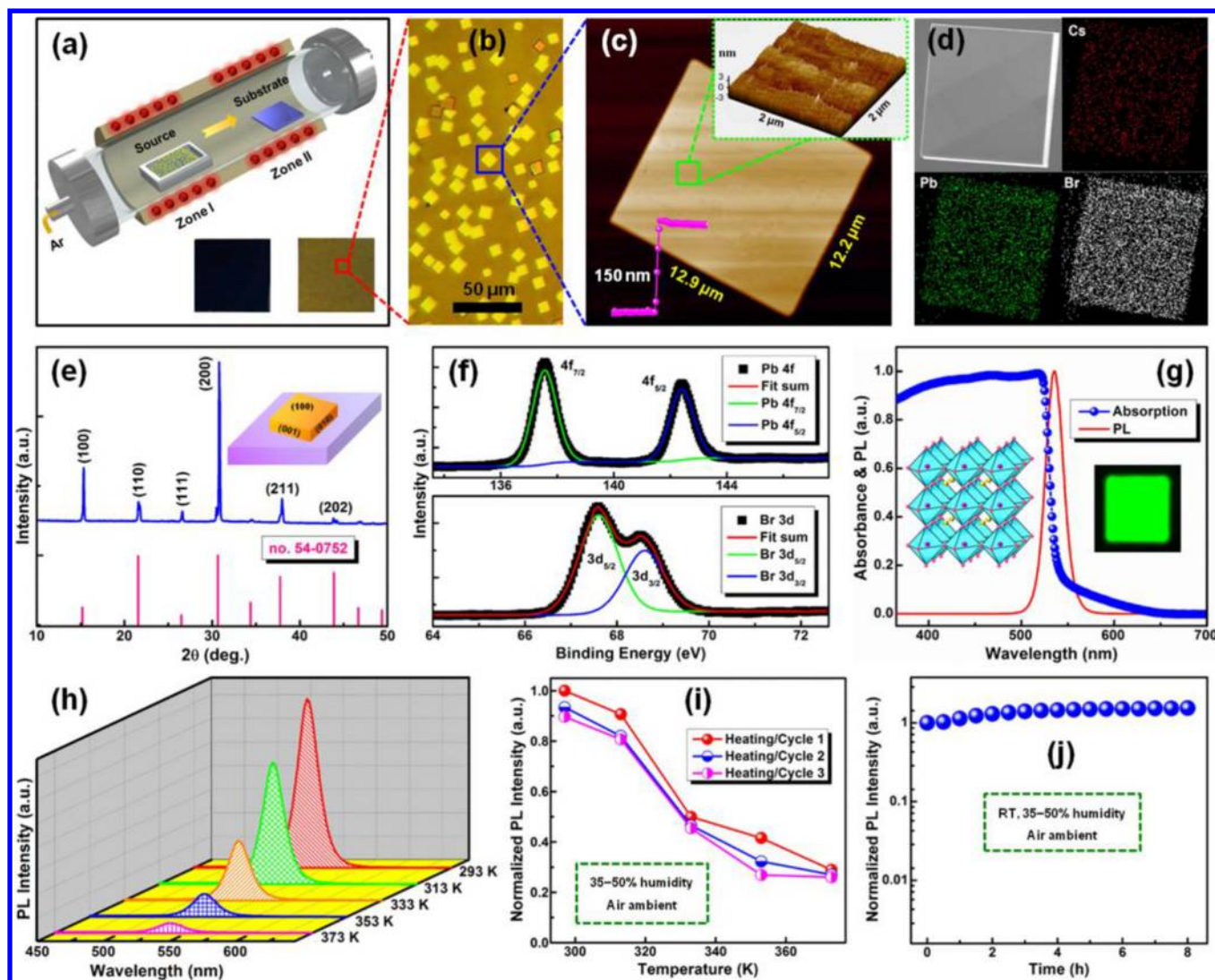


Figure 1. (a) Schematic diagram of the reaction setup for CsPbBr₃ microplatelets. (b) Optical microscopy image of the as-prepared CsPbBr₃ sample grown on SiO₂/Si substrate. (c) 2D AFM image of a single CsPbBr₃ microplatelet. (d) EDS mapping of an individual CsPbBr₃ microplatelet showing the uniform distribution of Cs, Pb, and Br elements. (e) XRD spectrum (top) of the CsPbBr₃ microplatelets. The bottom pattern shows the standard cubic CsPbBr₃ XRD pattern. (f) XPS spectra of Pb 4f and Br 3d of the CsPbBr₃ microplatelets. (g) PL and absorption spectra tested at RT. (h) Temperature-dependent emission spectra. (i) Three thermal cycling tests of the CsPbBr₃ microplatelets. (j) Photostability test of the CsPbBr₃ microplatelets.

and response speed were substantially enhanced by introducing the localized surface plasmonic effect.²⁸

Nevertheless, there are many obstacles preventing the further improvement of the detector performance, for example, the presence of a good deal of grain boundaries and defects in nano- or polycrystalline material forms.^{32,33} Thus, the photocurrent signal generated in the photodetectors has to preferentially compensate for such undesirable current loss, resulting in a slower photoresponse and a lower photoresponsivity. Moreover, in previous studies, the CsPbX₃ light absorbers are often characterized by the high-density and low-size nanocrystallites,¹² which are short of well-defined facets and irregularly stacked. The trap states may exist at the interfacial areas among the nanocrystallites. In this way, the photogenerated carriers can be rapidly trapped. Hence, from the device point of view, it is very essential to further improve the material quality and maintain the integrity of CsPbX₃ light absorbers. Besides, in previous studies, the photodetector performance including the photoresponsivity, on/off photo-

current ratio, response time, and other related parameters has always been emphasized, but the critical stability of the unencapsulated device has been rarely studied. For a broader range of applications in low-cost photodetection, investigation of the working stability of the perovskite photodetectors is imperative and should be facilitated.

In this work, we report a vapor-phase approach for the growth of high-quality micrometer-scale inorganic CsPbBr₃ microplatelets. A high crystallinity and fewer grain boundaries make the perovskites less susceptible to degradation, resulting in a device with long-term operation stability. Temperature-dependent steady-state and time-resolved photoluminescence (PL) tests were first performed to study the emission mechanisms and carrier recombination dynamics of the high-quality CsPbBr₃ microplatelets. Furthermore, a prototypical perovskite photodetector was prepared with a coplanar photoconduction configuration, exhibiting a high on/off ratio of 4.6×10^3 , a photoresponsivity of ~ 1.33 A/W, and a specific detectivity of 0.86×10^{12} jones. More importantly, the working

stability of the device was investigated to prove the sensitivity of the photodetection ability to atmosphere temperature and to evaluate its compatibility for practical applications. The underlying reasons for the photoresponsivity and on/off ratio decays are also discussed briefly. It is reasonably believed that our results may promote the development of stable and efficient perovskite photodetectors suitable for practical applications even under harsh environments.

RESULTS AND DISCUSSION

In this experiment, high-quality micrometer-scale inorganic CsPbBr₃ microplatelets were synthesized by using the chemical vapor phase deposition (CVD) technique, in which a double temperature zone tube furnace was employed as the reactor. The corresponding setup diagram is illustrated in Figure 1a, and the detailed synthesis process is shown in the Experimental Section. The inset in Figure 1a presents the photographs of the CsPbBr₃ microplatelets grown on the commercially available SiO₂/Si substrate. From the areas marked by a red square, we can see that the microplatelets are characterized by well-defined square shapes with smooth outer surfaces and sharp edges (Figure 1b). Note that the lateral dimension of the CsPbBr₃ microplatelets distributes in a range of 7–15 μm with a typical thickness of ~150 nm. Besides, the produced CsPbBr₃ microplatelets present a highly smooth surface, as shown in the two-dimensional (2D) atom force microscopy (AFM) image shown in Figure 1c, and the corresponding surface roughness is about 0.6 nm (inset in Figure 1c), which is perfectly flat in optical level and also is evidence of the highly crystallized nature of CsPbBr₃ microplatelets. Energy-dispersive X-ray spectroscopy (EDS) tests were used to study the elemental composition and distribution of an individual CsPbBr₃ microplatelet. As displayed in Figure 1d, the Cs, Pb, and Br elements are distributed uniformly across the whole microplatelet and overlay with each other perfectly. Quantitative analysis of the spectrum reveals that the Cs, Pb, and Br elements have an atomic ratio of about 1:1.1:3.2 (Figure S1, Supporting Information), which is close to the stoichiometry of the CsPbBr₃ material. In fact, many comparative experiments have been carried out to achieve such high-quality CsPbBr₃ microplatelets. Other forms were also prepared at relatively low reaction pressure, such as CsPbBr₃ thin films and microwire networks, as presented in Figure S2 (Supporting Information). The morphology dependence of CsPbBr₃ products on reaction pressure may be due to their different thermodynamically crystalline process in CsPbBr₃ formation. Furthermore, X-ray diffraction (XRD) measurements were used to examine the structural characteristics of CsPbBr₃ microplatelets. As shown in Figure 1e, two dominant and characteristic diffraction peaks at 15.34° and 30.81° were clearly observed, corresponding to the (100) and (200) crystalline planes of cubic CsPbBr₃ (*a* = 2.92 Å, JCPDS no. 54-0752). Note that the line width of the (100) diffraction peak is as small as 0.03°, which implies that the CsPbBr₃ microplatelets are well-fabricated with a high crystalline quality. The cubic-structured phase and single-crystal nature of the CsPbBr₃ microplatelets were further verified by the transmission electron microscopy (TEM) measurements (Figure S3, Supporting Information). The above structural investigation reveals that the microplatelets grow along the [010] and [001] directions and the (100) facet is exposed. We further carried out X-ray photoelectron spectroscopy (XPS) measurements to study the bond states and chemical composition of the as-

prepared CsPbBr₃ microplatelets. Figure S4 shows the total XPS spectrum of the products, where the signals from Br, Pb, Cs, and Br elements are observed above the detection limit. The Pb 4f and Br 3d signals as measured are displayed in Figure 1f. As shown in the upper pane, the binding energies for Pb 4f_{5/2} and Pb 4f_{7/2} are 139.8 and 135 eV, respectively, which are related to the stoichiometric CsPbBr₃, or Pb²⁺ in PbBr₃^{−34}. As shown in the bottom pane, a binding energy (65.0 eV) corresponding to Br 3d_{5/2} and a splitting corresponding to the Br 3d_{3/2} (1.0 eV) can be observed. This component can be assigned to the Br[−] state in CsPbBr₃.³⁵ Meanwhile, considering the atomic sensitivity factors and the area of the spectra, the atomic ratio of Pb and Br is approximated to be 1:3.04, matching well with the results in EDS measurements. To investigate the optical quality of the CsPbBr₃ microplatelets, PL and absorption spectra were both measured at room temperature (RT, 293 K). Figure 1g shows a highly symmetric PL spectrum peaking at 530 nm with a relatively narrow full-width at half-maximum (fwhm) of 27 nm. The blue line presents the absorption spectra of the CsPbBr₃ microplatelets, and a sharp absorption edge occurs at 540 nm, consistent with previous studies.^{14,31} The inset in Figure 1g illustrates a typical PL image of a single CsPbBr₃ microplatelet, and the green emission is uniformly distributed over the whole microplatelet. In addition, it is noteworthy that the undesirable sub-bandgap emissions caused by the defect states were not detected in the present case. The remarkable optical properties suggest that the CsPbBr₃ microplatelets have a low defect density, which is of great importance for their applications as a light absorber in photodetectors. In addition, the emission wavelength of the perovskite microplatelets can be adjusted with different halide composition. As displayed in Figure S5, by mixing the Cl/Br, the emission wavelength of CsPbBr_{3−x}Cl_x (0 < *x* < 3) was adjusted from blue to green emission. By mixing the Br/I, the emission wavelength of CsPbBr_{3−x}I_x (0 < *x* < 3) was adjusted from green to red emission.

To assess the suitability of the as-synthesized CsPbBr₃ microplatelets to practical photodetector applications, we performed the thermal cycling PL tests (293–373 K) for the CsPbBr₃ microplatelets. Also, the photostability measurements under a continuous light excitation (365 nm) were carried out. As shown in Figure 1h, a series of PL spectra at various ambient temperatures were measured under a constant light excitation and put together for comparison. One can see that the PL intensity decreases gradually with increasing the ambient temperature because of the added nonradiative recombination induced by the well-known heating effect. At 313 K, the emission intensity reduced to 75%, and at 373 K, the PL intensity remains ~25% of the initial value. Furthermore, we performed three successive thermal cycling tests to study the thermal stability. As shown in Figure 1i, an emission decay of only ~10% occurs after three heating–cooling cycles, greatly superior to the performance of CsPbBr₃ thin films (Figure S6) and other forms.^{36,37} To get more convincing conclusions regarding thermal stability of CsPbBr₃ microplatelets, 10 successive thermal cycling measurements were further conducted. As presented in Figure S7 (Supporting Information), only a ~20% integrated PL intensity decay was produced after 10 cycles. The results show the reproducible and stable features of the CsPbBr₃ light absorber and also its remarkable temperature resistance, opening opportunities for making temperature-stable photodetectors. Moreover, the photostability of the CsPbBr₃ microplatelets was investigated

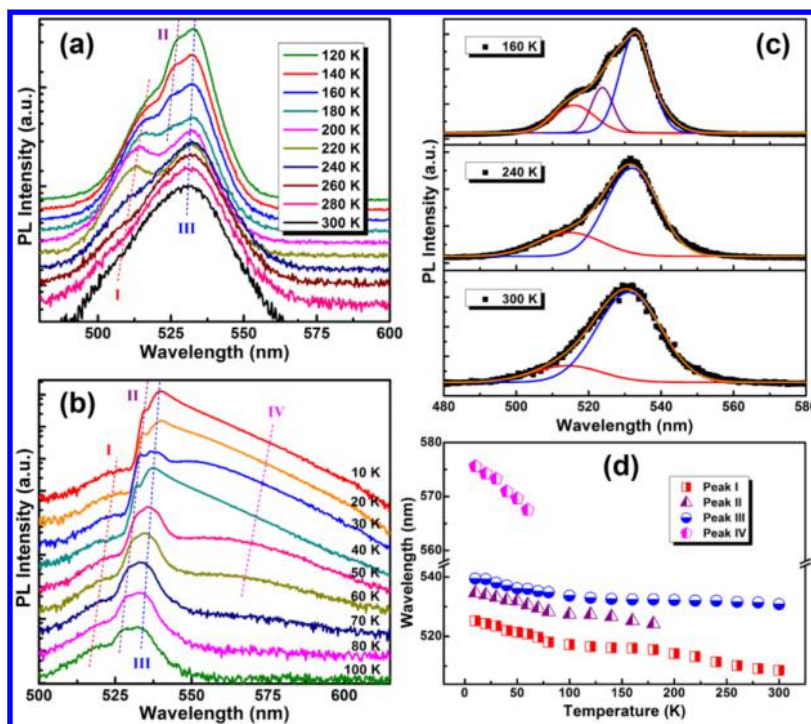


Figure 2. Temperature-dependent PL spectra ranging from (a) 300 to 120 K and (b) 100 to 10 K on a semilogarithmic scale. (c) Gaussian deconvolution measured at 160, 240, and 300 K. (d) Shift of the emission positions of peaks I, II, III, and IV as a function of measured temperature.

by illuminating the sample with a portable ultraviolet lamp at a 10 cm distance. Figure S8 displays a series of PL spectra of the CsPbBr₃ microplatelets recorded every 0.5 h. It is unexpectedly yet interesting that a PL enhancement behavior was observed over the test period (8 h), which is presumably related to the “photoactivation” phenomenon as reported in a previous study.³⁸

To better understand the optical transition mechanisms of the perovskite microplatelets, the PL measurements were further conducted in the temperature range of 10–300 K. As displayed in Figure 2a,b, the PL spectra at two temperature windows were put together for a clear view of the emission evolution process. At 300 K, one can observe that the PL spectrum is characterized by an intense emission peak at around 530 nm (peak III), accompanied by an imperceptible shoulder at ~505 nm (peak I). With a decrease of the temperature, the intensity ratio of the two peaks changes. Here, peak III was assigned to the free exciton emission in CsPbBr₃ microplatelets. Many reports have indicated that the exciton binding energy of CsPbBr₃ is larger than 26 meV, ensuring excitons survival well above RT, which is well consistent with our discussions.^{39,40} The relatively weak peak I, which locates at the high-energy side of peak III, may be related with the electron–light hole transition in CsPbBr₃, and the previous studies have also reported a similar phenomenon.^{41,42} Of course, more work should be done to further investigate the transition mechanisms of peak I. With the further decrease of temperature, the intensity of peaks I and III increases, and their peak positions gradually blue-shift. Herein, we consider that the electron–phonon interaction and the crystal lattice thermal expansion should be responsible for such an uncommon blue-shift phenomenon, and a similar observation has been reported in lead/copper chalcogenide semiconductors.⁴³ As the temperature decreased to 180 K, a new emission peak appeared at the position of 525 nm (peak II), and this emission component

can be retained at the low temperature of 10 K. Du and Shi have performed the density functional theory calculations to determine the origin of the above-bandgap emission in CsPbBr₃.⁴⁴ The energy level of Br vacancy (V_{Br}) locates above the conduction band of CsPbBr₃. Experimentally, Sebastian et al. observed a bound excitonic emission in CsPbBr₃ with the energy level located above the free exciton emission.⁴⁵ Considering the energy separation between peaks II and III (free excitons), we attributed peak II to V_{Br} -related emission in CsPbBr₃ microplatelets. It is noteworthy that another emission component at the long-wavelength side, ~565 nm (peak IV), appears when the temperature reduces to 60 K, and its intensity proportion increases gradually with decreasing temperature, which may result from the reduction of nonradiative recombination probability at low temperatures. Also, a similar phenomenon has been observed in other studies.^{46,47} Therefore, we attributed peak IV to the trapped charge-carrier recombination or trap-mediated radiative recombination. Because of the absence of irregular and/or reversal PL shift, we believe that there is no structural phase transition in the cooling process, unlike previous reports on CH₃NH₃PbI₃ systems.^{48,49}

The above results on temperature-dependent PL spectra suggest that there are four carrier transition channels in CsPbBr₃ microplatelets, and they take effect at different temperature levels. In the 200–300 K temperature window, peak I and III dominate the PL spectra, as displayed clearly by the Gaussian deconvolution of the emission spectra in Figure 2c (300 and 240 K, bottom and middle panes). In the 70–180 K temperature window, peaks I, II, and III contribute the emission spectra, as shown in a Gaussian deconvolution performed at 160 K (upper pane). In the 10–60 K temperature window, peaks I, II, III, and IV constitute the PL spectra together. In Figure 2d, the variation of peak positions for four emission contributors is summarized as a

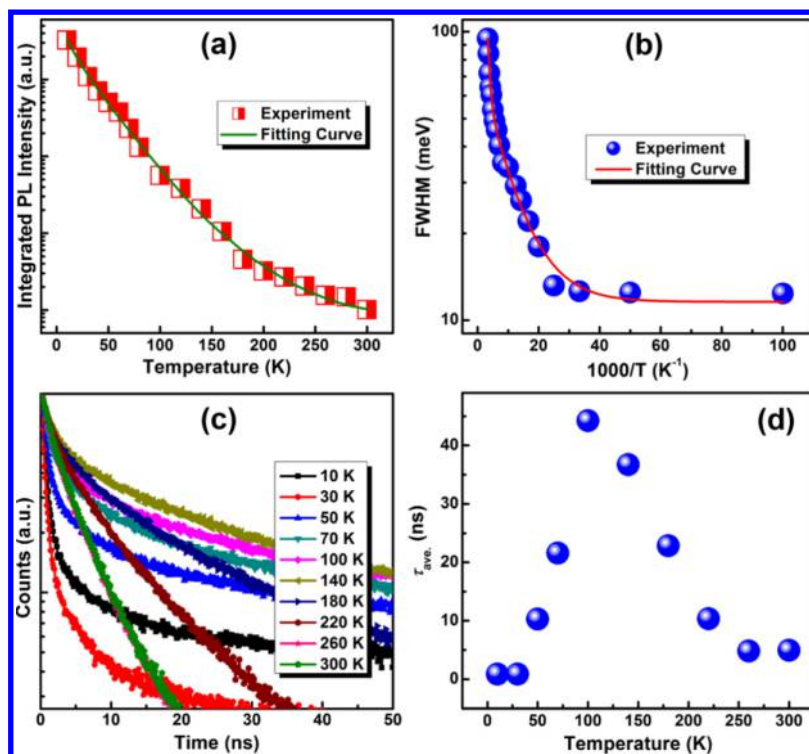


Figure 3. (a) Integrated PL intensity and (b) fwhm of the CsPbBr₃ microplatelets versus inverse temperature and the fitted lines. (c) PL decays at different temperature points. (d) Changing trend of τ_{ave} of the CsPbBr₃ microplatelets.

function of temperature. From the variation trend of peak III, the temperature sensitivity of the photon energy of CsPbBr₃ microplatelets can be estimated according to the empirical equation of $E_T = E_0 + \eta T$, in which E_T and E_0 are the measured photon energy and fitting constant at $T = 0$. Finally, the temperature sensitivity coefficient of $\eta = dE/dT$ was estimated to be 0.316 ± 0.007 meV/K.

Because the cubic phase of CsPbBr₃ microplatelets is stable in the temperature range, important parameters like exciton binding energy (E_b) and exciton–phonon interaction are estimated using the data procured above. As shown in Figure 3a, the intensity of peak III (exciton-related emission) as a function of temperature is plotted to determine the E_b . One can see that the integrated PL intensity increases exponentially with decreasing temperature owing to the thermally activated nonradiative recombination process, and eq 1 was used to describe such temperature quenching behavior⁵⁰

$$I(T) = \frac{I_0}{1 + A \exp\left(-\frac{E_b}{KT}\right)} \quad (1)$$

where A is a proportional constant, I_0 the PL intensity at 0 K, and K the Boltzmann constant. The value of E_b is 31.6 ± 1.3 meV, slightly larger than the RT thermal ionization energy. In addition, the fwhm of the PL emission of CsPbBr₃ microplatelets is fitted to investigate the exciton–phonon interaction in the carrier transition process as presented in Figure 3b, and the data are fitted using the equation⁵¹

$$\Gamma(T) = \Gamma_0 + \sigma T + \frac{\Gamma_{\text{op}}}{\exp(\hbar\omega_{\text{op}}/KT) - 1} \quad (2)$$

in which σ and Γ_{op} are the exciton–acoustic phonon interaction and the exciton–optical phonon contribution to

the line width broadening; $\hbar\omega_{\text{op}}$ and Γ_0 represent optical phonon energy and the inhomogeneous broadening contribution. It can be seen that Γ_0 is the main contributor to the formula at the low-temperature level, whereas σ and Γ_{op} dominate the formula at the high-temperature level. Thus, the value of fwhm would rise with temperature nonlinearly. Based on the fitting results, one conclusion that could be made is that the optical phonon is involved in the exciton–phonon interaction mainly. Accordingly, the values of $\hbar\omega_{\text{op}}$ and Γ_{op} were estimated to be 32.8 ± 1.9 and 66.9 ± 3.2 meV, respectively, implying a strong interaction between exciton and optical phonon and agreeing well with the thermal anti-quenching effect of CsPbBr₃ microplatelets at high temperature.

Furthermore, temperature-dependent transient PL tests were studied to understand the carrier recombination dynamics. Figure 3c displays a series of PL decay curves at 10 representative temperature points, which can be well-fitted with the biexponential decays⁵²

$$A(t) = A_1 \exp\left(-\frac{t}{\tau_1}\right) + A_2 \exp\left(-\frac{t}{\tau_2}\right) \quad (3)$$

in which A_1 and A_2 are the distribution coefficients and τ_1 and τ_2 stand for the characteristic PL decay lifetimes. The average lifetime (τ_{ave}) is evaluated with the A_i and τ_i values according to the expression of $\tau_{\text{ave}} = (A_1\tau_1^2 + A_2\tau_2^2)/(A_1\tau_1 + A_2\tau_2)$, and the obtained τ_{ave} values are summarized in Figure 3d. One can observe that the τ_{ave} increases gradually as the temperature increases below 100 K and then decreases significantly. We can infer from previous research that the lengthening of τ_{ave} at elevated temperatures may be due to the introduction of charge carrier trapping or localized states as the relaxation pathways.⁵³ Besides, the interaction between excitons and trap

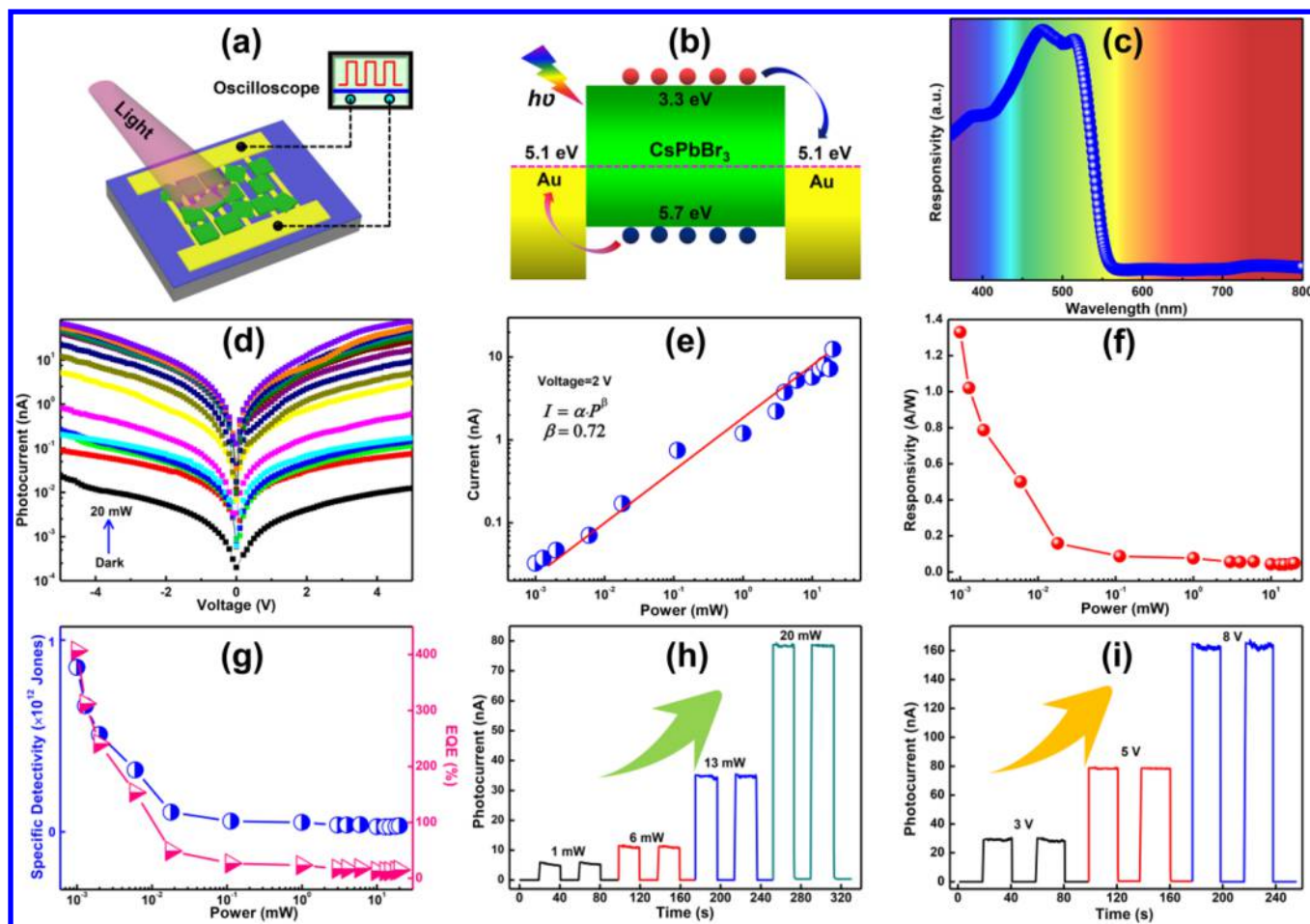


Figure 4. (a) Schematic diagram of the photodetector based on the CsPbBr₃ microplatelets. (b) Energy band alignment of the Au/CsPbBr₃/Au structure. (c) Spectral response of the CsPbBr₃ microplatelets photodetector. (d) *I*–*V* curves with different light irradiation intensities. (e) Logarithmic plot of the photocurrent versus light irradiation power at a bias of 2 V. (f) Photoresponsivity, (g) specific detectivity, and EQE versus light irradiation power. *I*–*t* curves operated at different (h) bias voltage and (i) light irradiation power.

states with rising temperature will induce the increase of PL lifetime. As for the subsequent decrease of τ_{ave} (100 K and above), the thermal degradation of CsPbBr₃ microplatelets at high temperatures deserves concern because the formation of added nonradiative recombination centers is inevitable.

To prove the potential applications of the CsPbBr₃ microplatelets in a photoelectric device, a photodetector device was prepared with a metal–semiconductor–metal configuration, and the corresponding device structure is shown in Figure 4a. Figure 4b shows the energy band diagram of the Au/CsPbBr₃/Au structure, in which E_C and E_V indicate the conduction band minimum and valence band maximum of the CsPbBr₃, respectively. When irradiated by the light source (405 nm in present case), electron–hole pairs are produced because of the photoelectric effect. Under the applied electric field, the carriers are separated rapidly and collected by opposite contact electrodes. Thus, the CsPbBr₃ microplatelets photodetectors could operate effectively, and accordingly, the photocurrent gradually increases with electric field intensity and light illumination. Figure 4c demonstrates the spectral response of the device, and a good spectral selectivity is well consistent with the absorption result shown above. Furthermore, the light intensity-dependent photoresponse performance of the studied photodetector was studied by changing the incident light intensity from the dark to 20 mW, and the

corresponding current–voltage (*I*–*V*) characteristics are depicted in Figure 4d. One can see that the photocurrent is highly dependent on the light intensity and increases gradually with increasing light intensity, which may be because the amount of photogenerated electron–hole pairs is proportional to the absorbed photon flux. Additionally, some *I*–*V* characteristic curves are featured by an asymmetrical behavior, which is probably because of the undesirable interface states at the Au/CsPbBr₃ interface. Note that the studied CsPbBr₃ microplatelets photodetector presents a low dark current of 1.73×10^{-11} A at 5.0 V, and the photocurrent increases several orders of magnitude to a value of 7.89×10^{-8} A at the illumination power of 20 mW. Thus, an on/off photocurrent ratio of 4.56×10^3 is achieved. In addition, the relationship between photocurrent (I_p) and the light excitation power was fitted by the power law of $I_p = \alpha P^\beta$, in which α represents a proportional constant, P is the light excitation power, and β is an exponent associated with its photoresponse corresponding to the light intensity. As shown in Figure 4e, the fitting curve produces a sublinear relationship of $\beta = 0.72$. This fractional power dependence ($0.5 < \beta < 1.0$) has also been found in other photodetectors and probably results from the complex process of electron–hole generation, recombination, and trapping in a semiconductor.²³ Besides, other key device parameters—photoresponsivity (*R*), specific detectivity (*D**)

Table 1. Summary of the Performance of the Perovskite-Based Photodetectors

active materials	responsivity (A/W)	D^* ($\times 10^9$ jones)	EQE (%)	on/off ratio	ref
$\text{CH}_3\text{NH}_3\text{PbI}_3$ single crystals	2.531	—	396.2	—	54
$\text{CH}_3\text{NH}_3\text{PbBr}_3$ single crystals	0.012	—	3.17	—	55
$\text{CH}_3\text{NH}_3\text{PbI}_3$ films	0.49×10^{-6}	—	—	70	56
$\text{CH}_3\text{NH}_3\text{PbBr}_{3-x}\text{I}_x$ films	0.055	—	—	$\sim 10^3$	57
$\text{CH}_3\text{NH}_3\text{PbI}_3$ nanowire	0.1	1020	—	340	58
$\text{CH}_3\text{NH}_3\text{PbI}_3$ microwire	1.32	2500	—	23	59
$\text{CH}_3\text{NH}_3\text{PbBr}_3$ nanowire	—	—	61.9	62	60
CsPbBr_3 nanosheets	0.64	—	54	$\sim 10^3$	30
CsPbBr_3 nanocrystals	0.01	1.684	8	$\sim 10^5$	61
$\text{CsPbBr}_{3-x}\text{I}_x$ nanocrystals	>1.0	—	—	$\sim 10^4$	62
CsPbBr_3 microcrystals	2.1	—	485	$\sim 10^3$	63
CsPbBr_3 microparticles	0.18	—	42	8.0×10^3	64
CsPbBr_3 microplatelets	1.33	1800	407	4.6×10^3	this work

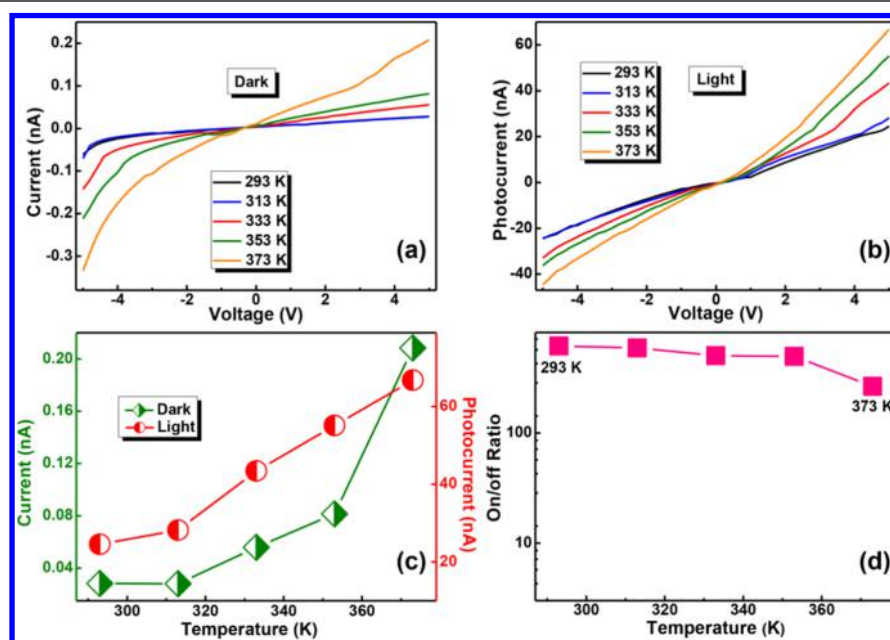


Figure 5. I – V characteristics of the photodetector (a) in dark conditions and (b) under light irradiance with different temperature conditions. (c) Variation of the dark current and photocurrent of the photodetector at different temperature points (5 V, 6 mW). (d) Variation of the on/off current ratio of the CsPbBr_3 microplatelets photodetector at different temperature points.

and external quantum efficiency (EQE)—were examined to quantitatively evaluate its photoresponse ability. R is defined by the equation $R = (I_p - I_d)/P_{\text{opt}} \cdot S$, where I_d and P_{opt} stand for the dark current and the incident light intensity, respectively, and S is the active area. Figure 4f presents the photoresponsivity curve versus the incident light power, and a maximum value of 1.33 A/W was achieved under a low light excitation power of 1 μW . The undesired saturation at high light illumination power may be related with the photocurrent saturation, which results from the filling mechanism of the sensitizing centers in CsPbBr_3 microplatelets at high light illumination intensity. As presented in Figure 4g, the D^* and EQE were also calculated and plotted with the incident light illumination power using the following relations:

$$D^* = \frac{(A \cdot f)^{1/2}}{\text{NEP}} \quad (4)$$

$$\text{NEP} = \frac{(\overline{i_n^2})^{1/2}}{R} \quad (5)$$

$$\text{EQE} = \frac{I_p/e}{P_{\text{opt}}/h\nu} = R \frac{hc}{e\lambda} \quad (6)$$

in which A represents the effective area of the sample; NEP is the noise equivalent power; f is the electrical bandwidth; $(\overline{i_n^2})^{1/2}$ is the root-mean-square value of the noise current acquired by using a lock-in amplifier at different frequencies; λ is the light wavelength, and h and c stand for Planck's constant and light speed. Figure S9 presents the relationship of $(\overline{i_n^2})^{1/2}$ with the frequency, from which the noise level per unit bandwidth (1 Hz) of the photodetector was measured to be 0.76 pA Hz^{1/2}. Then, the value of D^* was calculated and plotted with light illumination intensity. Both parameters display almost the same evolution trend as photoresponsivity, with the highest values of 0.86×10^{12} jones and 407% at the incident illumination power of 1 μW , which are much better than many reported results based on lead halide perovskites, as seen in Table 1.

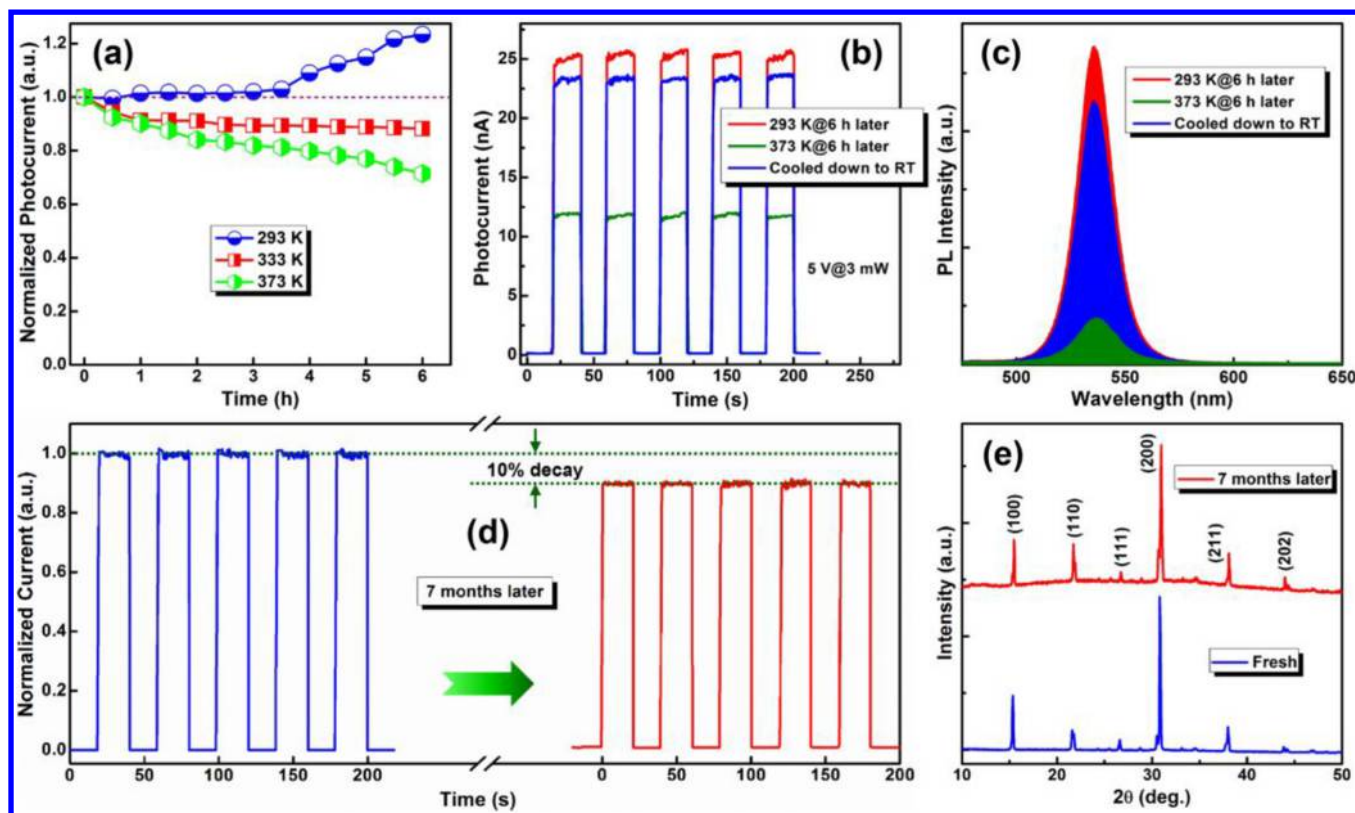


Figure 6. (a) Normalized photocurrent of the unencapsulated CsPbBr₃ microplatelets photodetector versus running time at 293, 333, and 373 K. (b) Temporal photoresponse curves of the device after aging for 6 h at 293 and 373 K and after naturally cooling to RT. (c) PL spectra of the CsPbBr₃ microplatelets after aging for 6 h at 293 and 373 K and after naturally cooling to RT. (d) Long-term storage stability. The last five cycles are the photoresponse after storage in air for 7 months. (e) Comparison of the XRD patterns of the CsPbBr₃ microplatelets before and after storage for 7 months in ambient air.

Temporal photoresponse is also an important performance parameter for the photodetectors. The current–time (I – t) curves of the device versus bias voltage were therefore measured to examine the time response, where the pulse bandwidth of the excitation light was fixed at 20 s per cycle and the illumination was switched on and off at different bias voltages (3, 5, and 8 V). As presented in Figure 4h, the current reversibly switches between low and high conductance with good stability and reproducibility, producing a series of tunable on/off photocurrent ratios. Furthermore, using the photocurrent response recorded at 8 V, the photoresponse time (rise time and fall time), which are the key evaluation criteria for detection performance, are obtained to be 20.9 and 24.6 ms, as displayed in Figure S10. As for the relatively slow response speed, two possible reasons should be considered. First, in our case, the CsPbBr₃ microplatelets were dispersed in toluene and then transferred to the SiO₂/Si substrate with Au patterned fingers. Therefore, inevitable interface defects may exist between CsPbBr₃ microplatelets and the Au electrode. Second, during the device testing, the incident light is radiated from the upper portion onto the photodetector, that is, on the top of the CsPbBr₃ microplatelets. Owing to the relatively large thickness of about 150 nm, the photogenerated carriers need more time to travel through the microplatelets to the underlying Au electrode, resulting in a relatively slow response time. In addition, the photoresponse with different light irradiance powers at 5 V was also studied. As displayed in Figure 4i, the photocurrent increases gradually with rising irradiance power, and the reproducible and prompt photocurrent response also

indicates a good cycling stability. In addition to the visible light detection, the photodetector is also sensitive to ultraviolet light. Figure S11 displays the corresponding I – V curves and time-dependent photoresponse of the device under 265 nm light illumination. Similarly, the current can reversibly switch between low and high conductance when the light irradiance was turned on and off, demonstrating a remarkable stability and reproducibility.

It is well-known that the working stability is a challenging issue for perovskite-based photodetectors.⁶⁵ In actual application, the photodetector has to be operated in the open air with inconsistent conditions. Therefore, the working stability for the studied photodetector without encapsulation was further investigated at different ambient temperatures to evaluate its temperature tolerance, and the corresponding I – V curves in dark conditions and under light illumination (405 nm, 6 mW) were measured. As shown in Figure 5a, one can observe that the dark current increases gradually with increasing ambient temperature, which can be well-fitted by the following Arrhenius equation:

$$A = A_0 \exp\left(-\frac{E_a}{KT}\right) \quad (7)$$

where A_0 is a proportional constant and K and T are the Boltzmann constant and temperature, respectively. The fitting curve yields the activation energy of $E_a \approx 0.27 \pm 0.08$ eV. For perovskites, the temperature-dependent current increase may be caused by the ion migration, generating halide vacancy defects. Mizusaki and co-workers have experimentally con-

firmed that the activation energy for the migration of Br^- in CsPbBr_3 is about 0.25 eV,⁶⁶ and the material becomes more conducting with the increase of temperature, well consistent with the results in our case. Of course, the thermally activated process for increased carriers in CsPbBr_3 microplatelets is an indisputable fact, similar to traditional semiconductors. Under the light illumination, the I - V curves display the same dependence on ambient temperature as that of the dark current. Typically, at a bias of 5 V, the photocurrent increases from 20 nA at RT to 62 nA at 373 K. Figure 5c summarizes the variation of dark current and photocurrent of the photodetector at different temperature points. As presented in Table S1, the increasing rate of dark current from 293 to 373 K is 6.83, while the decreasing rate of on/off ratio ($I_{\text{light}}/I_{\text{dark}}$) from 293 to 373 K is 2.33. Obviously, the decreasing rate of $I_{\text{light}}/I_{\text{dark}}$ is smaller than the increasing rate of dark current. If the reduction of $I_{\text{light}}/I_{\text{dark}}$ ratio is just due to the increase of dark current, the changing rate of the two parameters should be equal. Therefore, we consider that other positive factors may exist, which compensate for the $I_{\text{light}}/I_{\text{dark}}$ ratio. For instance, the elevated ambient temperature could provide energy for the separation of electrons and holes; the increased charge carrier density at high temperature could lower the effective energy barrier height. The above two factors favor an easier carrier transport and contribute to an enhancement in photoconductivity. To further investigate the temperature resistance of the proposed photodetector, temperature-dependent I - t curves of the device on a semilogarithmic scale were measured, in which the bias voltage and light illumination were fixed 5 V and 6 mW, respectively. As shown in Figure S12, the dark current and photocurrent increase simultaneously with temperature as observed in Figure 5c. Even at a high ambient temperature of 373 K, the CsPbBr_3 microplatelets photodetector can operate well, demonstrating a reproducible and stable response to on/off cycles. The results show that the photodetector has a remarkable temperature resistance, suggesting CsPbBr_3 microplatelets are potentially compatible with practical applications even in harsh environments. No study about the high-temperature working ability of perovskite-based photodetectors has been reported to date, making our study original and useful for progressing to a practical photodetector. Based on the obtained data from temperature-dependent I - t curves, the on/off ratio at different temperatures was calculated and is summarized in Figure 5d. One can see that the on/off photocurrent ratio exhibits a slight decay of $\sim 16\%$ at 353 K. As the ambient temperature is increased to 373 K, the on/off current ratio decreases by 43%. Despite this, the photodetection ability of such a photodetector can be efficiently sustained. We therefore attribute this decrease of on/off ratio to the increased defect states caused by the thermal effect, by which the photogenerated carriers are quickly trapped. In addition, the photoresponsivity and noise current of the photodetector at different temperature points were also evaluated. As shown in Figure S13, some distinct but regular changing trends are revealed. The photoresponsivity and noise current increase simultaneously with temperature.

To illuminate the long-term working stability of the studied photodetector, the aging test was performed at three representative temperature points (293, 333, and 373 K) under ambient air conditions. The photocurrent of the device without encapsulation versus the running time was performed at fixed bias voltage (5 V) and light excitation power (3 mW).

An interesting phenomenon was observed in Figure 6a: the photocurrent recorded at 293 K displays an increasing trend over the whole test period (6 h), which may be due to the "photoactivation" behavior as observed in photostability measurements of the CsPbBr_3 microplatelets. Presumably, the photoactivation behavior favors the removal of dangling bonds or undesirable surface states by highly energetic photons; thus, the transport ability of carriers could be improved to some extent. Of course, the heating effect at such a temperature level is insufficient to produce an increase of structural defects that benefit the trap states formation in perovskites. Correspondingly, the dark current of the device over a 6 h running time (Figure S14) does not display an obvious change except for a slight fluctuation. Therefore, the phenomenon of ion migration could be ruled out in our case. As the ambient temperature is elevated to 333 K, an undesired photocurrent decay of $\sim 10\%$ occurs. When the ambient temperature is further increased to 373 K, the decay magnitude of the photocurrent increases accordingly, and a decay value of $\sim 33\%$ is produced after 6 h of continuous operation. The corresponding temporal photoresponse curves of the device after aging for 6 h at 293 and 373 K are presented in Figure 6b, and the significantly reduced photocurrent is well consistent with the above discussion. It is understandable that the heating effect would increase with the ambient temperature and running time of the device. Still, the temperature resistance in our study is quite superior to that of other reports.^{67,68} An additional and very important observation is that such a photocurrent decay is recoverable; after cooling to RT, the photocurrent (blue line, Figure 6b) could recover to 92% of the initial level. The above studies prove the stable and recoverable performance of the CsPbBr_3 microplatelets photodetector even operated at high temperatures, motivating the existing applications and suggesting additional potential applications. To better understand the temperature resistance of CsPbBr_3 microplatelets, we recorded the PL spectra after aging for 6 h at 293 and 373 K and after being cooled to RT naturally. As displayed in Figure 6c, the effective PL emission can be sustained even after a 6 h aging test at 373 K except for a substantially reduced PL intensity compared with the counterpart recorded at 293 K (after 6 h of aging), which implies that the optical performance of the CsPbBr_3 microplatelets suffers from a degradation owing to the activation of added nonradiative pathways at higher temperature. However, what is exciting is that such a PL emission decay is just a transient behavior. After the sample is naturally cooled to RT, the PL performance recovers to $>90\%$ of its original value. Therefore, in terms of the material stability only, the CsPbBr_3 microplatelets demonstrate an obvious and impressive advantage compared with other perovskite material systems. Maybe this is the essential reason that ensuring the remarkable temperature resistance of the studied photodetector even worked in the open air and without encapsulation.

Moreover, the long-term storage stability of the present photodetector was studied. As presented in Figure 6d, the last five cycles correspond to the temporal photoresponse curve of the photodetector after seven months of storage in air ambient conditions, without any encapsulation and protection. One can see that the device can still work properly and maintains its photodetection ability except for a slight photocurrent decay of about 10%, obviously better than the devices constructed with hybrid perovskite photodetectors with poor material stability.⁶⁹ Figure 6e shows the XRD patterns of the CsPbBr_3 micro-

platelets before microplatelets before and after seven months of storage to demonstrate their environmental stability. It can be observed that all the diffraction peaks belong to the cubic CsPbBr₃, and no other diffraction signals were found except for the slight change of the peak intensity, suggesting that the CsPbBr₃ microplatelets have the ability to maintain the phase purity for seven months of storage. The good material stability together with the remarkable device performance imply that the CsPbBr₃ microplatelets are reliable and effective materials for photodetector applications.

CONCLUSION

We have successfully demonstrated a novel configuration of photodetectors by using vapor-processed micrometer-scale inorganic CsPbBr₃ microplatelets as the light absorber. Temperature-dependent and time-resolved PL measurements were first performed to study the emission mechanisms and carrier recombination dynamics of the CsPbBr₃ microplatelets. Photoresponse analysis reveals that the device possesses a stable and repeatable photocurrent response, and a high on/off ratio of 4.5×10^3 , a responsivity of 1.33 A/W, and a specific detectivity of 0.86×10^{12} Jones were achieved in the prepared CsPbBr₃ microplatelets photodetector. In addition, the unencapsulated photodetector demonstrates an excellent working stability against oxygen and water degradation over the temperature endurance test, and the photodetection ability can be efficiently sustained at 373 K, indicating a good temperature resistance and a desired compatibility for practical applications under harsh environments. The underlying reasons of remarkable working stability have been studied on the basis of thermal cycling and long-term PL measurements of CsPbBr₃ microplatelets. This work demonstrates that vapor-processed micrometer-scale CsPbBr₃ microplatelets are attractive candidates to be applied in stable and high-performance photodetectors.

EXPERIMENTAL SECTION

Synthesis of CsPbBr₃ Microplatelets. In the experiments, commercially available SiO₂/Si template was used as the substrate, and argon was used as the carrier gas. Particularly, a mixture of CsBr and PbBr₂ powders with a molar ratio of 1:2 was used as the precursors, which was heated in temperature zone I of the tube furnace at 580 °C. The precleaned SiO₂/Si substrate was located at the center of temperature zone II, and the corresponding temperature was set at 360 °C. After the furnace was heated to the preset temperature, the argon flow rate was set at 100 sccm, and the internal pressure was fixed at 180 Torr. The typical growth time for the CsPbBr₃ products was 20 min.

Device Preparation. The CsPbBr₃ microplatelets photodetector device was fabricated to investigate their optoelectronic applications. The CsPbBr₃ microplatelets dispersed in toluene were sonically agitated for 10 s and then drop-casted onto the SiO₂/Si substrate with Au patterned fingers (note that Au contact electrodes fingers were 5 μm wide and 360 μm long, and the channel spacing is 5 μm). Finally, the samples were annealed at 60 °C for 10 min on a hot plate to decrease the contact resistance and increase the contact adhesion.

Characterizations. *Materials.* Scanning electron microscopy (SEM) (Jeol-7500F, 15 keV) was used to study the microstructures. AFM (Dimension Icon, Bruker) was performed to research the thickness and surface roughness. XRD (Panalytical X'Pert Pro) was performed to study the crystallinity. The EDS attached SEM was used to investigate the element compositions of the products. The crystal structure of the as-synthesized CsPbBr₃ microplatelets was analyzed using TEM with a JEM-3010 electron microscope. XPS was performed with a SPECS XR50 system with Mg Kα (1254 eV,

width 0.7 eV) as the X-ray source. The absorption of the CsPbBr₃ microplatelets was measured using a Shimadzu UV-3150 spectrophotometer. The steady-state PL spectra were recorded by a double-grating spectrofluorometer (Horiba; Fluorolog-3) equipped with a digital temperature controller (LakeShore-325) and a closed-cycle helium cryostat (Jannis CCS-100). The time-resolved PL was tested by a lifetime measurement system with a pulsed LED (Horiba; 371 nm, pulse duration: 1.4 ns) as the light source. Optical images were captured by a Leica SP2 instrument.

Devices. The electrical and optoelectrical measurements of the device were performed by a system that combines a semiconductor characterization system (Keithley 4200-SCS), light sources, a monochromator, an oscilloscope (Tektronix, DPO2012B), and an optical chopper (SRS, SR540) in air ambient. For temperature-dependent optoelectrical measurements, the photodetector was installed on a ceramic heating plate, and the operation temperature could be adjusted from RT to 443 K with a precision of 0.5 K.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: [10.1021/acs.chemmater.8b02435](https://doi.org/10.1021/acs.chemmater.8b02435).

Additional experimental data as described in the text (PDF)

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Notes

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